

# INNER PLANETS OF SOLAR SYSTEM

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*For Carmen*

ABSTRACT. This paper provides a visualization of the four inner planets of the solar system.

## PREFACE

In this paper I visualize the *orbits* of the *inner* planets: *Mercury*, *Venus*, *Earth* and *Mars*. These orbits are *planar curves* in *space*. Since KEPLER (1571 – 1630) we know the shape of these curves: they are *ellipses*, which we need to put into their correct relative positions and orientation in different orbital planes in space.

The visualization is provided by figures drawn with METAPOST, where I disregard the inclination of the orbits towards the ecliptic. The presentation is kept at high school level with an emphasis of introducing ideas in a conceptual manner, though not losing sight of computation.

Berlin, 24 April 2010

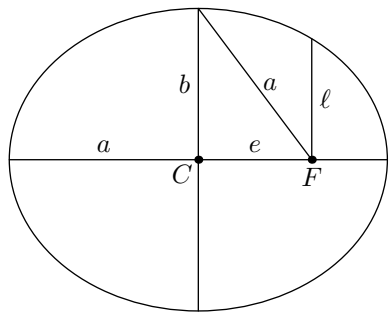
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## 1. ORBITS OF PLANETS

1.1. **Inside the orbital plane.** As a mathematical warm-up we will start with some planar geometry. We recall the *fundamental invariants* determining the shape of an ellipse, as in section *Conics* of COXETER [1, §8.4], or in more detail in [3, §1.2].

The fundamental invariants of ellipses are the *numerical eccentricity*  $\varepsilon$  and the *semilatus rectum*  $\ell$ . The shape (*ovalness*) is determined by  $\varepsilon$ , the size by  $\ell$ . In astronomical context rather than  $\ell$  the *semimajor axis*  $a$  is given.



The entities in the figure are related as follows

$$a = \frac{\ell}{1 - \varepsilon^2} \quad b = \frac{\ell}{\sqrt{1 - \varepsilon^2}}$$
$$e = \varepsilon \cdot a \quad a^2 = b^2 + e^2$$

$b$  is the *semiminor axis*,  $e$  is the *linear eccentricity* see [3, (1.9), (1.12), (1.14), (1.15)]

The equation of the ellipse in the coordinate system of the axes is given by [3, (1.13)]

$$\frac{x^2}{a^2} + \frac{y^2}{b^2} = 1$$

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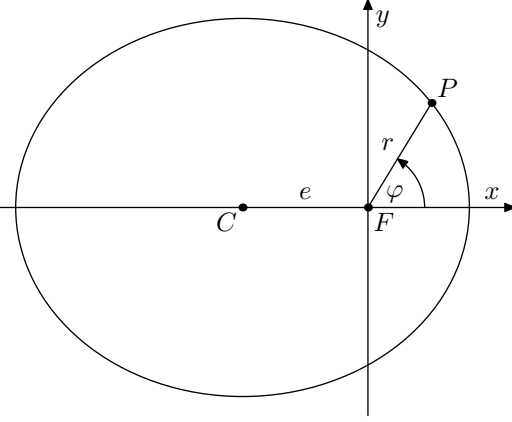
When this is written  $x^2 + (ay/b)^2 = a^2$ , the ellipse appears as a circle of radius  $a$  scaled down in the  $y$ -direction by the scaling factor  $b/a = \sqrt{1 - \varepsilon^2}$ .

When the focus  $F$  is the origin the ellipse equation reads

$$(1) \quad (x + e)^2 + \frac{y^2}{1 - \varepsilon^2} = a^2$$

From this one we derive the equation in polar coordinates

$$(2) \quad r = \frac{\ell}{1 + \varepsilon \cos \varphi}$$



*Proof.* By (1)

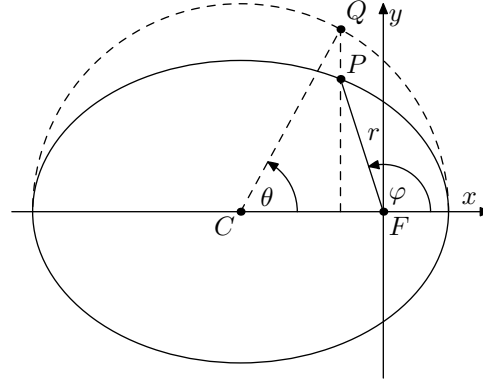
$$0 = (1 - \varepsilon^2)(x^2 + 2ex + e^2 - a^2) + y^2 = x^2 + y^2 - \varepsilon^2 x^2 + 2(1 - \varepsilon^2)ex - (1 - \varepsilon^2)b^2 =$$

$$= r^2 - \varepsilon^2 x^2 + 2\ell \varepsilon x - \ell^2 = r^2 - (\varepsilon x - \ell)^2 = (r + \varepsilon x - \ell)(r - \varepsilon x + \ell)$$

giving us two possible polar equations for the right focal point,  $r(1 + \varepsilon \cos \varphi) = \ell$  and  $r(\varepsilon \cos \varphi - 1) = \ell$ ; the latter only gives points for  $\varepsilon > 1$  (*hyperbola*) as  $r > 0$ .  $\square$

In the next figure we let the point  $P = (x, y) = (r \cos \varphi, r \sin \varphi)$  move on the ellipse and consider its twin  $Q = (x, ay/b)$  on the circle of radius  $a$  around  $C = (-e, 0)$

$$(3) \quad \begin{aligned} \sin \theta &= \frac{\sqrt{1 - \varepsilon^2} \sin \varphi}{1 + \varepsilon \cos \varphi} \\ \cos \theta &= \frac{\cos \varphi + \varepsilon}{1 + \varepsilon \cos \varphi} \\ \sin \varphi &= \frac{\sqrt{1 - \varepsilon^2} \sin \theta}{1 - \varepsilon \cos \theta} \\ \cos \varphi &= \frac{\cos \theta - \varepsilon}{1 - \varepsilon \cos \theta} \\ \tan \frac{\varphi}{2} &= \sqrt{\frac{1 + \varepsilon}{1 - \varepsilon}} \tan \frac{\theta}{2} \end{aligned}$$



Polar coordinates of  $Q$  imply  $a \sin \theta = ay/b$  and  $a \cos \theta = x + e$ , so by (2):

$$\begin{aligned} \sin \theta &= \frac{y}{b} = \frac{r \sin \varphi}{b} = \frac{\sqrt{1 - \varepsilon^2} \sin \varphi}{1 + \varepsilon \cos \varphi} \\ \cos \theta &= \frac{x + e}{a} = \frac{r \cos \varphi + \varepsilon a}{a} = \frac{(1 - \varepsilon^2) \cos \varphi}{1 + \varepsilon \cos \varphi} + \varepsilon = \\ &= \frac{\cos \varphi - \varepsilon^2 \cos \varphi + \varepsilon + \varepsilon^2 \cos \varphi}{1 + \varepsilon \cos \varphi} = \frac{\cos \varphi + \varepsilon}{1 + \varepsilon \cos \varphi} \end{aligned}$$

giving the first two relations in (3). The others are immediate as well, for example

$$\tan \frac{\theta}{2} = \frac{\sin \theta}{1 + \cos \theta} = \frac{\sqrt{1 - \varepsilon^2} \sin \varphi}{1 + \varepsilon \cos \varphi + \cos \varphi + \varepsilon} = \frac{\sqrt{1 - \varepsilon^2}}{1 + \varepsilon} \frac{\sin \varphi}{1 + \cos \varphi}$$

We also note the formula (directly implied by (3)):

$$(4) \quad r = a(1 - \varepsilon \cos \theta) = \frac{\ell}{1 + \varepsilon \cos \varphi}$$

The differential of (4) gives  $dr = a\varepsilon \sin\theta d\theta$ ; similarly from (2) we obtain

$$\begin{aligned} r\varepsilon \sin\varphi d\varphi &= (1 + \varepsilon \cos\varphi) dr = a\varepsilon(1 + \varepsilon \cos\varphi) \sin\theta d\theta = && \text{by (3)} \\ &= a\sqrt{1 - \varepsilon^2} \varepsilon \sin\varphi d\theta \end{aligned}$$

hence

$$(5) \quad (1 - \varepsilon \cos\theta) d\varphi = \sqrt{1 - \varepsilon^2} d\theta$$

**1.2. Celestial mechanics.** Let  $V$  designate our 3-dimensional vector space with the sun in the origin. Let the vector  $v \in V$  point to the planet under consideration. By NEWTON's *gravitational law* we have

$$(6) \quad \ddot{v} = -\frac{\alpha}{r^3} v, \quad r = |v|,$$

where  $\alpha = G \cdot (M + m)$ ,  $G$  is the *gravitational constant*,  $M$  is the mass of the sun and  $m$  that of the planet.

*Remark.* In fact, we have put the barycentre at rest, the planet is at position  $\frac{M}{M+m}v$  and the sun is at position  $-\frac{m}{M+m}v$ . For our illustrations this will make no real difference, as the barycentres of all inner planets lie deep inside the sun. Anyway, a point representing the sun in a figure is about five times larger in ratio to the real sun.

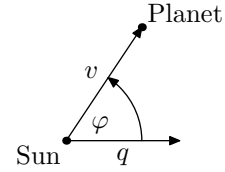
The vector  $c = v \times \dot{v}$  is constant:  $\dot{c} = v \times \ddot{v} = 0$ , as  $\ddot{v}$  is parallel to  $v$  by (6). Let  $L = \mathbf{R}c$  be the line spanned by  $c$  and  $P = L^\perp$  be the plane perpendicular to  $L$ , so  $V = L \oplus P$ . We have just seen that both  $v$  and  $\dot{v}$  lie in the plane  $P$  for all time:  $v, \dot{v} \in P$ , the planetary movement is *planar*.

Now we investigate the plane vector  $q = \dot{v} \times c - \frac{\alpha}{r} v$ . Differentiation gives

$$\dot{q} = \ddot{v} \times c - \frac{\alpha}{r} \dot{v} + \frac{\alpha \dot{r}}{r^2} v = -\frac{\alpha}{r^3} v \times c - \frac{\alpha}{r} \dot{v} + \frac{\alpha \dot{r}}{r^2} v = 0$$

as  $v \times c = v \times (v \times \dot{v}) = r\dot{v} - r^2\ddot{v}$  and  $v \cdot \dot{v} = r\dot{r}$  by deriving  $v \cdot v = r^2$ . So,  $q$  is an invariable vector in the orbital plane and we let  $\varphi$  be the angle between  $v$  and  $q$ , having thus introduced *polar coordinates*  $r = |v|$ ,  $\varphi = \angle(q, v)$  in  $P$ .

Now  $q \cdot v = v \cdot (\dot{v} \times c) - \alpha r = (v \times \dot{v}) \cdot c - \alpha r = c \cdot c - \alpha r$  or  $r(\alpha + |q| \cos\varphi) = |c|^2$ , which takes the form (2) with  $\varepsilon = |q|/\alpha$  and  $\ell = |c|^2/\alpha$ . This is KEPLER's first law.



In case  $\varepsilon < 1$  (ellipse) the minimal distance of the planet from sun  $r_{min} = \frac{\ell}{1+\varepsilon}$  is at  $\varphi = 0$  (*perihelion*) and the maximal distance  $r_{max} = \frac{\ell}{1-\varepsilon}$  is at  $\varphi = \pi$  (*aphelion*). The angle  $\varphi$  is also called the *true anomaly* and the angle  $\theta$  of the previous section 1.1 is called the *eccentric anomaly*.

Let us introduce an orthonormal frame  $(e_0, e_1, e_2)$  in  $V$  such that  $L = \mathbf{R}e_0$  and  $P = \mathbf{R}e_1 + \mathbf{R}e_2$ , with  $e_1$  pointing to the perihelion (so  $q = |q|e_1$ ) and  $e_2 = e_0 \times e_1$ . Then  $v = xe_1 + ye_2$  and  $\dot{v} = \dot{x}e_1 + \dot{y}e_2$ , hence  $c = v \times \dot{v} = (x\dot{y} - \dot{x}y)e_0$ . In polar coordinates we have

$$\begin{aligned} x &= r \cos\varphi & \dot{x} &= \dot{r} \cos\varphi - r\dot{\varphi} \sin\varphi \\ y &= r \sin\varphi & \dot{y} &= \dot{r} \sin\varphi + r\dot{\varphi} \cos\varphi \end{aligned}$$

and  $|c| = x\dot{y} - \dot{x}y = r^2\dot{\varphi}$ . This is the constancy of the *orbital angular momentum* of KEPLER's second law:

$$(7) \quad \sqrt{\ell\alpha} dt = r^2 d\varphi$$

The movement of planets on ellipses is not uniform, as  $r$  is not constant for  $\varepsilon > 0$ . The correct relation is the famous KEPLER equation (8) relating time and position. The proof also gives us KEPLER's third law (9).

$$(8) \quad \theta - \varepsilon \sin \theta = \frac{2\pi}{T}(t - t_0)$$

$$(9) \quad \frac{a^3}{T^2} = \frac{\alpha}{4\pi^2} = \frac{G(M + m)}{4\pi^2}$$

where  $T$  is the revolution time of the planet around the sun.

*Proof.* Substituting (4) in (7)  $\sqrt{\ell\alpha} dt = a^2(1 - \varepsilon \cos \theta)^2 d\varphi = a^2\sqrt{1 - \varepsilon^2}(1 - \varepsilon \cos \theta) d\theta$  by (5). We get by integration (where  $t_0$  corresponds to the time at perihelion):

$$\sqrt{\frac{\alpha}{a^3}} \int_{t_0}^t dt = \sqrt{\frac{\alpha}{a^3}}(t - t_0) = \int_0^\theta (1 - \varepsilon \cos \theta) d\theta = \theta - \varepsilon \sin \theta$$

After one revolution  $\theta = 2\pi$  and  $t - t_0 = T$ , completing the proof.  $\square$

## 2. ORBITAL ELEMENTS

The orbital plane of the Earth is called ecliptic, the position of the sun will be the origin and the  $x$ -axis points to the vernal point  $\Uparrow$ . I disregard the inclination of the other planets against the ecliptic.

For our purpose the approximate positions in STANDISH [4] are sufficient. The calculations of the elements, in particular for the *eccentric anomaly*  $\theta$  is done in the program `solar.c`, see appendix A.

The orbits are drawn as rotated ellipses with METAPOST with the sun in the origin. The values for the *true anomaly* are taken from the output of the `solar` program. The following visualisation of the orbits of the inner planets at their location on October 13, 2013 is generated by the METAPOST program:

```
% -----
%   Abbildung 5 -- inner planets
% -----

beginfig(5);
def orbit(expr ra,rb,an)=
  (fullcircle xscaled 2ra yscaled 2rb shifted -(ra+-+rb,0) rotated an)
enddef;
path vernal,part;
part = ((-0.25,0.25){2up+left}..(-0.2,0.5){right}.. (0,0){down});
vernal = part & reverse part reflectedabout ((0,0),up);
v=.4mm;
picture sun,mercury,venus,earth,mars;

draw fullcircle scaled 10v;
fill fullcircle scaled 2v;
sun = currentpicture;
clearit;

draw fullcircle scaled 5v shifted (0,.5v);
draw halfcircle scaled 5v rotated 180 shifted (0,5.5v);
draw (-2v,-4v)--(2v,-4v);
draw (0,-2v)--(0,-6v);
mercury = currentpicture;
clearit;
```

```

draw fullcircle scaled 5v shifted (0,.5v);
draw (-2v,-4v)--(2v,-4v);
draw (0,-2v)--(0,-6v);
venus = currentpicture;
clearit;

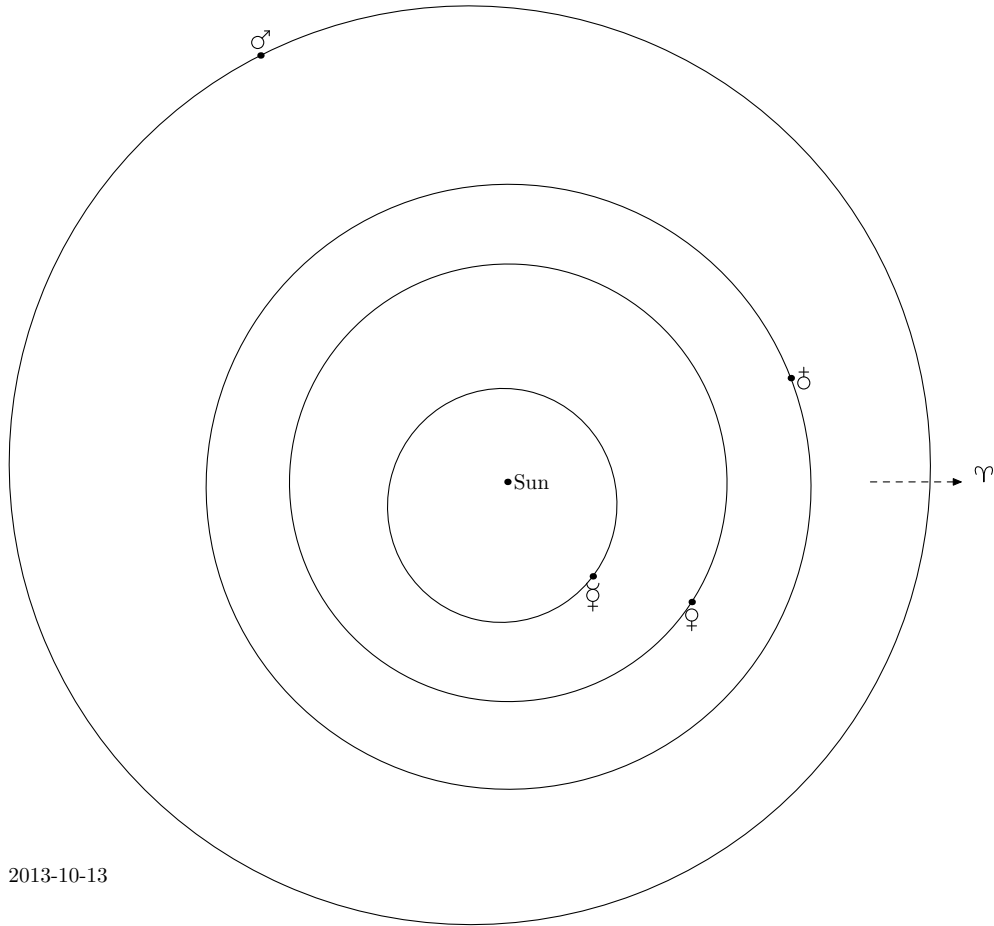
draw fullcircle scaled 5v shifted (0,-2.5v);
draw (-2v,2v)--(2v,2v);
draw (0,0)--(0,4v);
earth = currentpicture;
clearit;

path q[];
draw fullcircle scaled 5v shifted (0,-2.5v);
z10=(2.5/sqrt(2)*v,2.5/sqrt(2)*v);
z20=z10+(3v,3v);
q0=z10--z20;
q1=(x20-2v,y20)--z20;
q2=(x20,y20-2v)--z20;
draw q0 shifted (0,-2.5v);
draw q1 shifted (0,-2.5v);
draw q2 shifted (0,-2.5v);
mars = currentpicture;
clearit;

path p[];
z0=(0,0);
a1=3.871u; b1=0.9786*a1;
a2=7.233u; b2=0.99998*a2;
a3=10u; b3=0.99986*a3;
a4=15.237u; b4=0.996*a4;
% Sun
% Mercury
% Venus
% Earth
% Mars
% *** Output of solar 2013-10-13: ***
%julian : 2456579
%Mercury: -47.89, Perihel: 77.48
%Venus : -33.10, Perihel: 131.78
%Earth : 20.12, Perihel: 102.97
%Mars : 120.08, Perihel: -23.86
label.lft(btex 2013-10-13 etex,(-13u,-13u));
v1=-47.89; v2=-33.10; v3=20.12; v4=120.08;
w1=77.48; w2=131.78; w3= 102.97; w4=-23.86;
for i=1 upto 4:
  p[i]=orbit(a[i],b[i],w[i]);
  draw p[i];
endfor
for i=1 upto 4:
  z[i]=2a[i]*dir v[i];
  p[4+i]=z0--z[i];
  z[4+i]=p[i] intersectionpoint p[4+i];
endfor
dotlabel.bot(mercury,z5);
dotlabel.bot(venus,z6);
dotlabel.rt(earth,z7);
dotlabel.top(mars,z8);
draw sun;
z9=(15u,0);
drawarrow .8z9--z9 dashed evenly; % direction to vernal point

```

```
draw vernal scaled u shifted 1.05z9;
endfig;
% -----
```



#### APPENDIX A. THE PROGRAM `SOLAR.C`

The program `solar.c` calculates the true anomaly of the inner planets from a date. The calculation of the *julian* day number has been taken from the article *On calendar formulas* [2]. The routine `theta` solves the KEPLER equation (8).

```
//
```

```
/* -----*
```

```
Module:      solar.c
```

```
Description:
```

```
Calculating the true anomaly of inner planets
```

```
For details: http://berndt-schwerdtfeger.de/sol/solar.pdf
```

```
Input:  yyyy-mm-dd (-3000 <= yyyy <= +3000)
```

```
Output: direction of planets Mercury, Venus, Earth, Mars
```

Subroutines: theta, leap, qu

-----  
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-----\*/

```
#include <stdio.h>
#include <stdlib.h>
#include <math.h>

// set constants - Keplerian elements from Standish

char *planet[] = {"Mercury", "Venus ", "Earth ", "Mars  "};
int em[13] = {0,0,31,59,90,120,151,181,212,243,273,304,334};
double e[4] = {0.20563661, 0.00676399, 0.01673163, 0.09336511};
double de[4]= {0.00002123,-0.00005107,-0.00003661, 0.00009149};
double L[4] = {252.25166724, 181.97970850, 100.46691572, -4.56813164};
double dL[4]= {149472.67486623, 58517.81560260, 35999.37306329, 19140.29934243};
double w[4] = {77.45771895, 131.76755713, 102.93005885, -23.91744784};
double dw[4]= { 0.15940013, 0.05679648, 0.31795260, 0.45223625};
#define PI 3.14159265358979324
#define DEG (180.0/PI)
#define RAD (PI/180.0)

// prototype statements for functions

double theta(double,double); // theta - eps*sin theta = m
int leap(long); // = 1 if year is a leap year
long qu(int,int); // = [a/b] Gaussklammer

// -----
// main program
// -----

int main(int argc, char *argv[]){

    if (argc == 1 || *argv[1] == '?'){

        printf("\nsolar v1.0, 2010-05-03 \n");
        printf("-- Copyright (C) 2010 Berndt E. Schwerdtfeger -- \n");
        printf("\nTrue anomaly of Mercury, Venus, Earth and Mars \n");
        printf(" Input:      yyyy-mm-dd (yyyy in range -3000 - 3000) \n");
```

```

    return EXIT_SUCCESS;           // end the program
}                                  // argc > 1 here

int i,d,m,y;
long j,x;
double T, c[4], f[4], l[4], t[4], v[4];

sscanf(argv[1], "%d-%d-%d", &y, &m, &d);

if (argc > 2) {
    printf("Too many parameters !\n");
    return EXIT_FAILURE;          // end the program
}
if (y < -3000 || y > 3000) {
    printf("Error: outside allowed range -3000 <= year <= +3000 !\n\n");
    return EXIT_FAILURE;          // end the program
}
if (y == 1582 && m == 10 && d > 4 && d < 15) {
    printf("This date is invalid, it did never exist !\n");
    return EXIT_FAILURE;          // end the program
}

if (y < 1582 || y == 1582 && m*100+d < 1005) // Julian era
    j = 1721423 + d + em[m] + (y-1)*365 + qu(y-1,4);
else {                                     // Gregorian era
    x = y - 2001;
    j = 2451910 + d + em[m] + x*365 + qu(x,4) - qu(x,100) + qu(x,400);
}
if (m > 2) j += leap(y);                  // correction for leap year

x = j - 2451545;                           // days since 2000-01-01
T = (double)x / (double)36525;              // julian centuries

for (i=0; i<4; i++) {
    c[i] = e[i] + de[i]*T;                  // eccentricity
    l[i] = L[i] + dL[i]*T;                  // mean longitude
    v[i] = w[i] + dw[i]*T;                  // longitude of perihelion
    f[i] = (l[i] - v[i])*RAD;               // mean anomaly
    t[i] = theta(f[i], c[i]);               // eccentric anomaly
    f[i] = 2*atan(sqrt((1+c[i])/(1-c[i]))*tan(t[i]/2))*DEG+v[i]; // true a.
}

printf("julian : %ld\n", j);
for (i=0; i<4; i++) printf("%s: %.2f, Perihel: %.2f\n", planet[i], f[i], v[i]);

return EXIT_SUCCESS;                       // exit the program
}

// -----
// subroutine: theta, leap, qu
// -----

double theta(double m, double e){          // solve Kepler equation
    double t, d, dt;

    while (m > PI) m -= 2*PI;                // normalize m ..
    while (m < -PI) m += 2*PI;              // .. to range -PI <= m <= PI

```



```

t = m + e*sin(m);

do
{ d = m - (t - e*sin(t));
  dt= d/(1-e*cos(t));
  t += dt;
} while (fabs(dt) > 0.0001);

return t;
}

// -----
int leap(long y){

int i = 0;
if (y%4 == 0)
  i = 1;
if (y > 1582 && y%100 == 0 && y%400 != 0)
  i = 0;
return i;
}

// -----
long qu(int a, int b){

long x = a/b;
if (a < 0) // for negative numerator ..
  x -= (a%b != 0); // .. if remainder, subtract 1
return x;
}
//

```

## REFERENCES

- [1] Harold Scott MacDonald Coxeter, *Introduction to Geometry*, 2nd ed., John Wiley & Sons, 1969.
- [2] Berndt E. Schwerdtfeger, *On calendar formulas* (2010), available at <http://berndt-schwerdtfeger.de/wp-content/uploads/pdf/calj.pdf>.
- [3] ———, *Invariants of Curves of second order* (2013), available at <http://berndt-schwerdtfeger.de/wp-content/uploads/pdf/c2.pdf>.
- [4] Erland Myles Standish, *Keplerian Elements for Approximate Positions of the Major Planets*, available at [http://ssd.jpl.nasa.gov/txt/aprx\\_pos\\_planets.pdf](http://ssd.jpl.nasa.gov/txt/aprx_pos_planets.pdf).